

## Influence of hot pressure on the magnetoresistance of CrO<sub>2</sub>

Xianjie Wang, Yu Sui, Xiudan Song, Ruibin Zhu, Zhengnan Qian, Wenhui Su, and Jinke Tang

Citation: *Journal of Applied Physics* **101**, 09J509 (2007); doi: 10.1063/1.2714272

View online: <http://dx.doi.org/10.1063/1.2714272>

View Table of Contents: <http://scitation.aip.org/content/aip/journal/jap/101/9?ver=pdfcov>

Published by the [AIP Publishing](#)

---

### Articles you may be interested in

[Magnetoresistance in granular CrO<sub>2</sub>: Effects of variation in crystallographic and magnetic microstructure](#)  
J. Appl. Phys. **101**, 103911 (2007); 10.1063/1.2733621

[Magnetic inhomogeneity and valence state in Sr<sub>2</sub>CrWO<sub>6</sub> double perovskite](#)  
J. Appl. Phys. **93**, 471 (2003); 10.1063/1.1524710

[Intergrain tunneling and magnetotransport properties in CrO<sub>2</sub> – RuO<sub>2</sub> / TiO<sub>2</sub>](#)  
J. Appl. Phys. **91**, 7923 (2002); 10.1063/1.1451882

[Extrinsic and intrinsic magnetoresistance contributions of CrO<sub>2</sub> thin films](#)  
J. Appl. Phys. **89**, 7699 (2001); 10.1063/1.1362658

[Characterization of the natural barriers of intergranular tunnel junctions: Cr<sub>2</sub>O<sub>3</sub> surface layers on CrO<sub>2</sub> nanoparticles](#)  
Appl. Phys. Lett. **77**, 2840 (2000); 10.1063/1.1320845

---



## Influence of hot pressure on the magnetoresistance of CrO<sub>2</sub>

Xianjie Wang

*Center for Condensed Matter Science and Technology (CCMST), Department of Physics, Harbin Institute of Technology, Harbin 150001, People's Republic of China*

Yu Sui<sup>a)</sup>

*Center for Condensed Matter Science and Technology (CCMST), Department of Physics, Harbin Institute of Technology, Harbin 150001, People's Republic of China and International Centre for Materials Physics, Chinese Academy of Sciences, Shenyang 110016, People's Republic of China*

Xiudan Song, Ruibin Zhu, and Zhengnan Qian

*Center for Condensed Matter Science and Technology (CCMST), Department of Physics, Harbin Institute of Technology, Harbin 150001, People's Republic of China*

Wenhui Su

*Center for Condensed Matter Science and Technology (CCMST), Department of Physics, Harbin Institute of Technology, Harbin 150001, People's Republic of China and International Centre for Materials Physics, Chinese Academy of Sciences, Shenyang 110016, People's Republic of China*

Jinke Tang

*Department of Physics, University of New Orleans, New Orleans, Louisiana 70148*

(Presented on 11 January 2007; received 31 October 2006; accepted 29 January 2007; published online 7 May 2007)

In this paper, we investigate the influence of high temperature and high pressure (hot pressure) on the magnetic and transport properties of polycrystalline CrO<sub>2</sub> samples compacted under high pressure and high temperature of up to 5 GPa and 600 °C, respectively. The magnetic moment increases with compacting temperature, and a metal-semiconductor transition is observed in hot-pressed samples, different from the cold-pressed samples. These results indicate that the formation of Cr<sub>2</sub>O<sub>3</sub> at the grain boundaries of CrO<sub>2</sub> is suppressed by hot pressure. The magnitude of low field magnetoresistance of up to 1 T at 5 K is enhanced first with the increase of compacting temperature and then decreased under higher compacting temperature. This result can be well explained by the change of spin-dependent tunneling at the modulated grain boundaries of CrO<sub>2</sub> due to the transformation from Cr<sub>2</sub>O<sub>3</sub> to CrO<sub>2</sub> under hot pressure. © 2007 American Institute of Physics. [DOI: 10.1063/1.2714272]

### I. INTRODUCTION

As the simplest half-metallic oxide predicted by band structure calculation,<sup>1</sup> the magnetotransport properties of CrO<sub>2</sub> have drawn much attention because it has nearly perfect spin polarization (close to 100%) and higher Curie temperature  $T_C$  of 395 K, which are necessary for spintronic device applications.<sup>2,3</sup> The large low field magnetoresistance (MR) originating from the spin-dependent carrier tunneling at grain boundary has been observed in CrO<sub>2</sub> granular samples at low temperature,<sup>3-9</sup> but only a very small value of MR can be obtained in CrO<sub>2</sub> at high temperature.

Usually, there is a naturally grown Cr<sub>2</sub>O<sub>3</sub> layer on the surface of CrO<sub>2</sub> particles,<sup>3,7,10</sup> in which the higher-order hopping of a carrier is considered to be one of the main reasons for the suppressed low field MR at high temperature.<sup>9</sup> It is obviously important to remove the Cr<sub>2</sub>O<sub>3</sub> layer for understanding the intrinsic magnetotransport property of CrO<sub>2</sub>. Unfortunately, Cr<sub>2</sub>O<sub>3</sub> is more stable than CrO<sub>2</sub> at atmospheric pressure so it is hard to eliminate this native Cr<sub>2</sub>O<sub>3</sub>

layer from the surface of CrO<sub>2</sub> particles in air. It is well known that CrO<sub>2</sub> is stable under high pressure. In addition, a high-pressure chamber can offer an isolated environment for CrO<sub>2</sub> to avoid losing oxygen and transforming into Cr<sub>2</sub>O<sub>3</sub>.<sup>11-13</sup> Therefore, it is worthwhile to modulate the Cr<sub>2</sub>O<sub>3</sub> layer by high temperature and high pressure (hot pressure), but clear experimental data on the influence of hot pressure on the magnetotransport property of CrO<sub>2</sub> are still lacking.

In this paper, we report some preliminary results of the magnetic and transport properties of CrO<sub>2</sub> samples compacted under hot pressure. Hot pressure not only prevents CrO<sub>2</sub> from transforming into Cr<sub>2</sub>O<sub>3</sub> but also changes the original Cr<sub>2</sub>O<sub>3</sub> to CrO<sub>2</sub>. An enhanced low field MR and the conductance of polycrystalline CrO<sub>2</sub> can be obtained under suitable hot pressure conditions.

### II. EXPERIMENTS

Nanosized CrO<sub>2</sub> powders were supplied by Dupont. They are needle-shaped particles with length of about 400 nm and an aspect ratio of about 9:1. These CrO<sub>2</sub> powders were pressed into a tablet and put into the chamber of a belt-type high-pressure apparatus, and compacted into bulk

<sup>a)</sup>Author to whom correspondence should be addressed; electronic mail: [suiyu@hit.edu.cn](mailto:suiyu@hit.edu.cn)

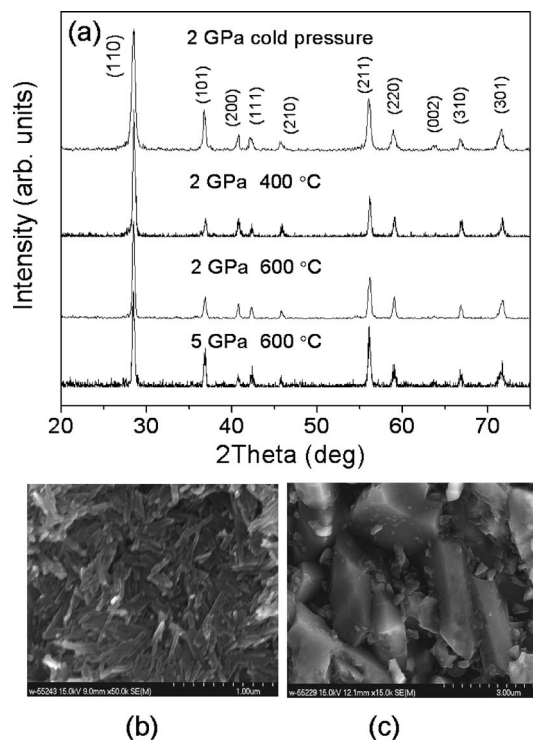


FIG. 1. (a) Room-temperature XRD patterns of CrO<sub>2</sub>, (b) SEM of 2 GPa cold-pressed sample, and (c) SEM of the 2 GPa, 500 °C sample.

samples under different pressures and temperatures of up to 5 GPa and 600 °C, respectively. In the hot pressure procedure, the pressure was first increased to a desired value, followed by temperature. After holding the pressure and temperature for 30 min, the temperature was decreased to room temperature first, and then the pressure was released. The measurements of the pressure and temperature are described in Ref. 14.

X-ray diffraction (XRD) patterns were collected using a Bede D<sup>1</sup> XRD spectrometer with Ni-filtered Cu K $\alpha$  radiation. The surface morphology of the samples was obtained using a Hitachi S-4700 field emission scanning electron microscope. The magnetic and transport properties were measured by using the physical properties measurement system (PPMS) of Quantum Design.

### III. RESULTS AND DISCUSSION

Figure 1(a) gives the room temperature XRD patterns of CrO<sub>2</sub> samples annealed under different pressures. Obviously, Cr<sub>2</sub>O<sub>3</sub> may be introduced by annealing the CrO<sub>2</sub> in air at 400 °C for 20 min,<sup>3,9</sup> but there is not peak of Cr<sub>2</sub>O<sub>3</sub> in all hot-pressed CrO<sub>2</sub> samples, indicating that high pressure can suppress the formation of Cr<sub>2</sub>O<sub>3</sub> up to 600 °C. Figures 1(b) and 1(c) show the scanning electron microscopy (SEM) micrographs for the cold-pressed and hot-pressed CrO<sub>2</sub>. The cold-pressed sample still has needle-shaped nanosized particles, but only large CrO<sub>2</sub> grains are seen in the hot-pressed sample due to the growth of CrO<sub>2</sub> particles under hot pressure. These results are consistent with the changes in the XRD patterns such that the peaks of annealed CrO<sub>2</sub> under high pressure are sharper than the cold-pressed samples.

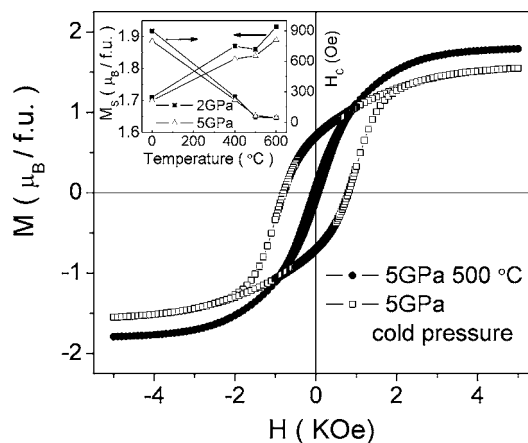


FIG. 2. Isothermal magnetization curves of CrO<sub>2</sub> at 5 K. Inset shows the variation of the  $M_S$  at 5 K and 5 T (left axes) and  $H_C$  (right axes).

Figure 2 shows the isothermal magnetization curves of CrO<sub>2</sub> at 5 K. The saturation moments per formula unit ( $M_S$ ) and coercivity ( $H_C$ ) of the cold-pressed sample are about  $1.7\mu_B$  and 900 Oe, respectively, similar to the original CrO<sub>2</sub> powders.<sup>15</sup> However, the  $H_C$  of the hot-pressed CrO<sub>2</sub> at 5 K decreases rapidly from 900 to 40 Oe with increasing compacting temperature from room temperature to 600 °C, as shown in the inset of Fig. 2. This feature is mainly caused by the increase of grain size due to the growth of CrO<sub>2</sub> particles under hot pressure and is consistent with the results of XRD and SEM. The  $M_S$  value of the cold-pressed CrO<sub>2</sub> ( $1.7\mu_B$ ) at 5 K and 5 T is much lower than the theoretical value of pure CrO<sub>2</sub> ( $2\mu_B$ ), as reported in literatures. It means that there is a large amount of Cr<sub>2</sub>O<sub>3</sub> on the surface of CrO<sub>2</sub> nanoparticles.<sup>16</sup> But it is obvious from the inset of Fig. 2 that  $M_S$  increases with increasing compacting temperature, and the highest value of  $M_S$  ( $1.93\mu_B$ ), close to the theoretical value, can be obtained at 2 GPa and 600 °C. This remarkable increase of  $M_S$  should be attributed to the transformation from antiferromagnetic Cr<sub>2</sub>O<sub>3</sub> to ferromagnetic CrO<sub>2</sub> on the grain surfaces. Because these data were measured in a high magnetic field of 5 T, the contribution due to the grain growth itself can be ignored. CrO<sub>2</sub> is more stable than Cr<sub>2</sub>O<sub>3</sub> under high pressure, thus the Cr<sub>2</sub>O<sub>3</sub> layer on the surface of the CrO<sub>2</sub> particles is eventually eliminated when they grow into larger particles under hot pressure. It should also be noticed that the values of  $M_S$  of 5 GPa hot-pressed samples are smaller than those at 2 GPa for the same compacting temperature, which suggests that the transformation from Cr<sub>2</sub>O<sub>3</sub> to CrO<sub>2</sub> is difficult under higher pressure because the amount of oxygen in the high-pressure chamber, which is necessary for oxidizing Cr<sub>2</sub>O<sub>3</sub> to form CrO<sub>2</sub>, decreases with increasing pressure. This result also gives strong evidence that the increase in  $M_S$  arises from the transformation of Cr<sub>2</sub>O<sub>3</sub> to CrO<sub>2</sub>. If it were the change of grain size that is mainly responsible for the increase in  $M_S$ , the  $M_S$  of the 5 GPa samples would be larger than the 2 GPa samples at the same compacting temperature because the former should have larger particles than the latter.

The temperature dependence of the resistivity ( $\rho$ ) of CrO<sub>2</sub> in zero field is displayed in Figs. 3(a) and 3(b). Here,  $\rho$

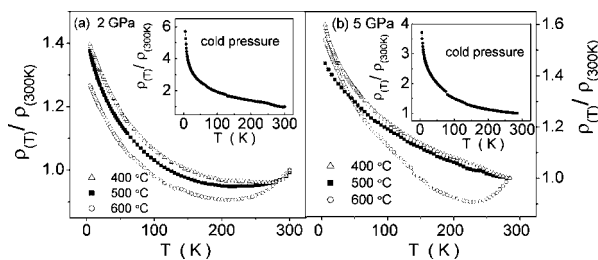


FIG. 3. Temperature dependence of the resistivity ( $\rho$ ) at zero fields of hot-pressed  $\text{CrO}_2$ : (a) 2 GPa sample and (b) 5 GPa sample. The insets of (a) and (b) are the cold-pressed samples of 2 and 5 GPa, respectively.

is normalized to the value at 300 K. It is clear that the  $\rho$  vs  $T$  curves of all cold-pressed samples show a semiconductor behavior in the whole temperature range, which is similar to that of original  $\text{CrO}_2$  powders.<sup>9</sup> However, the metal-semiconductor ( $M-I$ ) transition occurs in hot-pressed samples, indicating that the resistance at the grain boundaries is weakened because much of the original intergranular barriers, such as the  $\text{Cr}_2\text{O}_3$  layers, have been removed under hot pressure. It is interesting that the  $M-I$  transition can be found in all 2 GPa hot-pressed samples, but only observed in the 5 GPa hot-pressed sample at compacting temperature of 600 °C. We note that all samples showing the transition behavior have higher  $M_S$  over  $1.85\mu_B$ , and the 5 GPa samples have smaller  $M_S$  than the 2 GPa samples. It therefore suggests a strong correlation between the amount of  $\text{Cr}_2\text{O}_3$  at the grain boundaries and the electrical conductance of the samples. Under a pressure of 2 GPa, there is enough oxygen in the high-pressure chamber for the transformation from  $\text{Cr}_2\text{O}_3$  to  $\text{CrO}_2$  to occur. The transition occurs at temperatures as low as 400 °C and is easier at higher temperatures. The  $M-I$  transition is observed in all hot-pressed samples at 2 GPa. Under higher pressure of 5 GPa, the transition from  $\text{Cr}_2\text{O}_3$  to  $\text{CrO}_2$  becomes more difficult and much higher temperature is needed to achieve the same. Therefore, the  $M-I$  transition only occurs in the 5 GPa hot-pressed sample at 600 °C.

Figure 4 shows the MR of 5 GPa samples as a function of magnetic field at 5 K, and the inset shows the variation of MR at 1 T and 5 K for all samples. Here, MR is defined as  $\text{MR}(\%) = [\rho_H - \rho_{\max}] / \rho_H$ , where  $\rho_H$  and  $\rho_{\max}$  are the resistivity of the sample in external field  $H$  and the highest resistivity value, respectively. With increasing compacting temperature, the low field MR increases remarkably first and then decreases, as shown in the inset of Fig. 4. It is well known that the low field MR of granular  $\text{CrO}_2$  originates from the spin-dependent tunneling across the  $\text{Cr}_2\text{O}_3$  layer at the grain boundaries.<sup>3-8</sup> As discussed earlier, because the amount of  $\text{Cr}_2\text{O}_3$  in the 5 GPa hot-pressed samples is much more than that of the 2 GPa hot-pressed samples, the value of MR of the former is larger than that of the latter at the same compacting temperature. However, under lower compacting temperature, the  $\text{Cr}_2\text{O}_3$  layer as tunneling barrier still exists at the grain boundaries in the hot-pressed samples, which may become more suitable for spin-dependent tunneling under optimum hot pressure conditions. The largest MR value of 42% is observed in the sample at 5 GPa and 400 °C. With

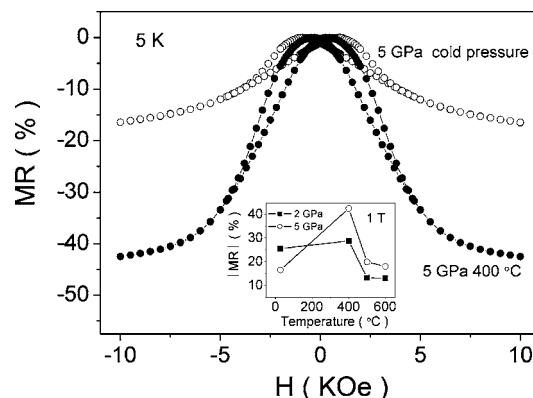


FIG. 4. The magnetic field dependence of MR at 5 K for the  $\text{CrO}_2$  samples; the inset shows the MR data at 1 T and 5 K.

increasing compacting temperature, the amount of the  $\text{Cr}_2\text{O}_3$  layer is reduced by the high pressure, and the low field MR is suppressed accordingly.

#### IV. CONCLUSIONS

In a high-pressure cell, not only the formation of  $\text{Cr}_2\text{O}_3$  at high temperature is suppressed, but the original  $\text{Cr}_2\text{O}_3$  layer on the  $\text{CrO}_2$  particle surfaces has also been transformed to  $\text{CrO}_2$ . The change in the amount of  $\text{Cr}_2\text{O}_3$  at the grain boundaries leads to the increase of magnetic moment and the appearance of the  $M-I$  transition in the hot-pressed  $\text{CrO}_2$  samples, as well as changes in the magnetotransport properties of  $\text{CrO}_2$ . As a result, the low field MR is remarkably enhanced by the hot pressure, and a large MR value of 42% was observed in the sample hot-pressed at 5 GPa and 400 °C.

#### ACKNOWLEDGMENT

This work was supported by the National Natural Science Foundation of China (Grant Nos. 10304004 and 50672019).

- <sup>1</sup>K. Schwarz, J. Phys. F: Met. Phys. **16**, L211 (1986).
- <sup>2</sup>T. Y. Cai, S. Ju, and Z. Y. Li, Appl. Phys. Lett. **88**, 192503 (2006).
- <sup>3</sup>J. B. Dai and J. K. Tang, Appl. Phys. Lett. **77**, 2840 (2000).
- <sup>4</sup>S. S. Manoharan, D. Elefant, G. Reiss, and J. B. Goodenough, Appl. Phys. Lett. **72**, 984 (1998).
- <sup>5</sup>J. M. D. Coey, A. E. Berkowitz, L. Balcells, and F. F. Putris, Phys. Rev. Lett. **80**, 3815 (1998).
- <sup>6</sup>K. Suzuki and P. M. Tedrow, Appl. Phys. Lett. **74**, 428 (1999).
- <sup>7</sup>J. P. Wang, P. Che, J. Feng, M. F. Lu, J. F. Liu, J. Men, Y. J. Hong, and J. K. Tang, J. Appl. Phys. **97**, 073907 (2005).
- <sup>8</sup>J. M. D. Coey, J. Appl. Phys. **85**, 5576 (1999).
- <sup>9</sup>J. B. Dai and J. K. Tang, Phys. Rev. B **63**, 064410 (2001).
- <sup>10</sup>A. Barry, J. M. D. Coey, and M. Viret, J. Phys.: Condens. Matter **12**, L173 (2000).
- <sup>11</sup>V. A. Sidorov, A. V. Rakhmanina, and O. A. Morya, Solid State Commun. **139**, 360 (2006).
- <sup>12</sup>S. F. Matar and G. Demazeau, Chem. Phys. Lett. **407**, 516 (2005).
- <sup>13</sup>B. Maddox, C.-S. Yoo, V. Iota, D. Kasinathan, W. Pickett, and R. Scalettar, Bull. Am. Phys. Soc. **49**, 408 (2004).
- <sup>14</sup>W. H. Su, D. M. Wu, X. Y. Li, X. F. Ma, J. S. Zhou, Z. N. Qian, Y. F. Wang, and W. N. Liu, Physica B **139-140**, 658 (1986).
- <sup>15</sup>J. B. Dai and J. K. Tang, Phys. Rev. B **63**, 054434 (2001).
- <sup>16</sup>N. A. Frey, S. Srinath, H. Srikanth, M. Varela, S. Pennycook, G. X. Miao, and A. Gupta, Phys. Rev. B **74**, 024420 (2006).